

**Riemannian Geometry, Differential Operators on Riemannian Manifolds**  
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## 1 Differential operators

In this notes, we always assume that  $M$  is a compact oriented Riemannian manifold with the Levi-Civita connection  $\nabla$ , and  $X, Y, \dots$  are smooth vector fields.

- $(p, q)$ -tensor is a smooth section of the bundle  $T_q^p M = (\otimes^p TM) \otimes (\otimes^q T^*M)$ .

- Contractions  $c_{ij} : T_q^p M \rightarrow T_{q-1}^{p-1}$  defined by

$$c_{ij}(x_1 \otimes \dots \otimes x_p \otimes y_1^* \otimes \dots \otimes y_q^*) = y_j^*(x_i) x_1 \otimes \dots \otimes \hat{x}_i \otimes \dots \otimes x_p \otimes y_1^* \otimes \dots \otimes \hat{y}_j^* \otimes \dots \otimes y_q^*.$$

- The interior product  $i_X$  of a vector field  $X$  with a covariant tensor of type  $(0, p)$  is the tensor  $c_{1,1}(X \otimes S)$  (of the type  $(0, p-1)$ ), defined as

$$(i_X(S))_m(x_1, \dots, x_{p-1}) = S_m(X(m), x_1, \dots, x_{p-1}).$$

- Extension of the covariant derivative: Let  $X$  be a vector field on  $X$ . The endomorphism  $D_X$  of  $\Gamma(TM)$  has a unique extension of the endomorphism of the space of tensors, still denoted by  $D_X$  which is type-preserving and satisfies the following

(1) For any  $S \in \Gamma(T_q^p)$  ( $p > 0, q > 0$ ) and any contraction  $c$ ,  $D_X(c(S)) = c(D_X S)$ .

(2) For any tensors  $S$  and  $T$ ,

$$D_X(S \otimes T) = D_X S \otimes T + S \otimes D_X T.$$

If for example  $S \in \Gamma(T_q^0 M)$ , the  $X_i$  are vector fields on  $M$ , then

$$(D_X S)(X_1, \dots, X_q) = X(S(X_1, \dots, X_q)) - \sum_{i=1}^q S(X_1, \dots, X_{i-1}, D_X X_i, \dots, X_q).$$

The readers should aware that, in general,

$$(D_X S)(X_1, \dots, X_n) \neq D_X(S(X_1, \dots, X_n)).$$

- Hessian: If  $f \in C^\infty(M)$ , then the  $(0, 2)$ -tensor  $Ddf$  is *symmetric*. The tensor  $Ddf$  is called the *Hessian* of  $f$ .

- Musical isomorphism:  $\omega(Y) = g(X, Y)$  gives the map between  $X$  and  $\omega$ . Flat  $\flat$ : the map  $X \rightarrow \omega$ ,  $\sharp$  is the inverse map.
- $div(X) := -tr(DX)$  and  $\delta\alpha = -trD\alpha$  (the co-differential of  $\alpha$ ) where  $DX$  is the endomorphism  $u \rightarrow D_u X$  and the  $D\alpha$  is a covariant two-tensor, where the trace is computed with respect to the Riemannian metric. Using the "musical isomorphism", these two notations can be viewed as equivalent. More general, for any  $(0, q)$ -tensor, we define the  $div(T)$  as a  $(0, q - 1)$  type tensor as  $\delta T = -tr_{12}DT$  where the notation  $_{12}$  just means that the trace is taken with respect to the first two variables. It is easy to check

$$\delta(fT) = f\delta T - i_{\nabla f}T.$$

- The Laplacian of  $f$  is the function  $\Delta f$  given by

$$\Delta f = div \nabla f = \delta df.$$

The Euclidean Laplacian on  $\mathbf{R}^n$  is given by

$$\Delta = -\sum_{i=1}^n \frac{\partial^2}{\partial x_i^2}.$$

- **Bochner's formula:** For any smooth functions on  $(M, g)$ , we have

$$g(\Delta(df), df) = |Ddf|^2 + \frac{1}{2} \Delta(|df|^2) + Ric(\nabla f, \nabla f).$$

- **The divergence operator.** Let  $X \in \Gamma(TM)$ . Define

$$div(X) = tr\{Y \rightarrow \nabla_Y X\}.$$

$div(X)$  is a smooth function on  $M$ . It is called the *divergence of  $X$* . The map  $div : \Gamma(TM) \rightarrow C^\infty(M)$  given by  $X \mapsto div(X)$  is called the *divergence operator*. In terms of local coordinate  $(U; x^i)$ , write

$$X|_U = X^i \frac{\partial}{\partial x^i},$$

then

$$\operatorname{div}(X) = \frac{1}{\sqrt{G}} \sum_{i=1}^m \frac{\partial}{\partial x^i} (\sqrt{G} X^i),$$

where  $G = \det(g_{ij})$ ,  $g_{ij} = g(\partial/\partial x^i, \partial/\partial x^j)$ .

*Proof.* Let  $(U, x^i)$  be a local coordinate, and write

$$X|_U = X^i \frac{\partial}{\partial x^i}.$$

Then

$$\nabla X = \left( \frac{\partial X^i}{\partial x^j} + X^k \Gamma_{kj}^i \right) dx^j \otimes \frac{\partial}{\partial x^i}.$$

Hence

$$\operatorname{div}(X) = \sum_{i=1}^m \left( \frac{\partial X^i}{\partial x^i} + X^k \Gamma_{ki}^i \right).$$

So we see that the div operator is a **differential operator of first order** acting on  $X$ .

By the formula,

$$\Gamma_{ki}^i = \frac{1}{2} g^{ij} \frac{\partial g_{ij}}{\partial x^k} = \frac{1}{2G} \frac{\partial G}{\partial x^k} = \frac{1}{\sqrt{G}} \frac{\partial \sqrt{G}}{\partial x^k}.$$

Thus

$$\operatorname{div}(X) = \frac{\partial X^i}{\partial x^i} + \frac{X^k}{\sqrt{G}} \frac{\partial \sqrt{G}}{\partial x^k} = \frac{1}{\sqrt{G}} \frac{\partial}{\partial x^i} (\sqrt{G} X^i).$$

- **The gradient of  $f$ .** Let  $f \in C^\infty(M)$ , define a tangent vector field  $\operatorname{grad}(f) \in \Gamma(TM)$ , by

$$g(\operatorname{grad}(f), X) = df(X) = X(f),$$

for every smooth tangent vector field  $X$ . The tangent vector field  $\operatorname{grad}(f)$  is called the *gradient* of  $f$ . In terms of local coordinate  $(U; x^i)$ ,

$$\operatorname{grad}(f) = \sum_{j=1}^m \left( \sum_{i=1}^m g^{ij} \frac{\partial f}{\partial x^i} \right) \frac{\partial}{\partial x^j},$$

where  $g_{ij} = g(\partial/\partial x^i, \partial/\partial x^j)$ ,  $(g^{ij}) = (g_{ij})^{-1}$ . Again, the “grad operator”  $\operatorname{grad} : C^\infty(M) \rightarrow \Gamma(TM)$  defined by  $f \mapsto \operatorname{grad}(f)$  is a **differential operator of first order** acting on  $f$ .

- **Beltrami-Laplace operator.** Let  $f \in C^\infty(M)$ , define  $\Delta f = \operatorname{div}(\operatorname{grad}(f))$ . It is called the *Beltrami-Laplace operator*. The operator  $\Delta f = \operatorname{div} \circ \operatorname{grad} : C^\infty(M) \rightarrow C^\infty(M)$  is a very important differential operator.

In local coordinate  $(U; x^i)$ ,

$$\Delta f = \frac{1}{\sqrt{G}} \sum_{i=1}^m \frac{\partial}{\partial x^i} \left( \sum_{j=1}^m \sqrt{G} g^{ij} \frac{\partial f}{\partial x^j} \right),$$

where  $G = \det(g_{ij})$ ,  $g_{ij} = g(\partial/\partial x^i, \partial/\partial x^j)$ ,  $(g^{ij}) = (g_{ij})^{-1}$ .

- **Volume form, interior product.** In a local coordinate  $(U; x^i)$ , let

$$\eta = \sqrt{G} dx^1 \wedge \cdots \wedge dx^m.$$

$\eta$  in fact is a global  $m$ -form, called *the volume form of  $M$* . Fix a smooth tangent vector field  $X$ , the interior product  $i(X)$  is defined by, for every tangent vector fields  $X_1, \dots, X_{m-1}$ ,

$$(i(X)\eta)(X_1, \dots, X_{m-1}) = \eta(X, X_1, \dots, X_{m-1}).$$

Then we have, for every smooth tangent vector field  $X$ ,

$$d(i(X)\eta) = \operatorname{div}(X)\eta.$$

*Proof:* By definition,  $\eta = \sqrt{G}dx^1 \wedge \dots \wedge dx^m$ . Hence, from the formula above,

$$\begin{aligned} \operatorname{div}(X)\eta &= \frac{\partial}{\partial x^i} (\sqrt{G}X^i) dx^1 \wedge \dots \wedge dx^m \\ &= \sum_{i=1}^m d((-1)^{i+1}\sqrt{G}X^i) \wedge dx^1 \wedge \dots \wedge \hat{dx}^i \wedge \dots \wedge dx^m. \end{aligned}$$

Write

$$\omega = \sum_{i=1}^m (-1)^{i+1}\sqrt{G}X^i dx^1 \wedge \dots \wedge \hat{dx}^i \wedge \dots \wedge dx^m.$$

It is easy to verify that  $\omega$  is independent of the choice of coordinates. So  $\omega$  is a globally defined  $(m-1)$ -form on  $M$ . The above identity gives

$$\operatorname{div}(X)\eta = d\omega.$$

It is easy to verify that  $\omega = i(X)\eta$  by definition.

- **The divergence theorem.** Use the above identity that  $\operatorname{div}(X)\eta = d\omega$ . and Stokes theorem, we get the divergence theorem: *Let  $(M, g)$  be a compact oriented Riemannian manifold, then, for every smooth tangent vector field  $X$ ,*

$$\int (\operatorname{div}X)\eta = 0,$$

where  $\eta$  is the volume form.

- **Global inner product for differential forms** We first define the inner product for differential forms. Let  $\phi, \psi$  are two  $r$ -forms. Let  $(U, x^i)$  be a local coordinate. We write

$$\phi|_U = \frac{1}{r!} \phi_{i_1 \dots i_r} dx^{i_1} \wedge \dots \wedge dx^{i_r},$$

$$\psi|_U = \frac{1}{r!} \psi_{j_1 \dots j_r} dx^{j_1} \wedge \dots \wedge dx^{j_r}.$$

We define, the inner product  $\langle \cdot, \cdot \rangle$  of  $\phi, \psi$  as

$$\langle \phi, \psi \rangle = \frac{1}{r!} \phi^{i_1 \dots i_r} \psi_{i_1 \dots i_r} = \sum_{i_1 < \dots < i_r} \phi^{i_1 \dots i_r} \psi_{i_1 \dots i_r},$$

where  $\phi^{i_1 \dots i_r} = g^{i_1 j_1} \dots g^{i_r j_r} \phi_{j_1 \dots j_r}$ . It is important to note that the definition is independent of the choice of local coordinates. We also have  $\langle \phi, \phi \rangle \geq 0$  and  $\langle \phi, \phi \rangle = 0$  if and only if  $\phi = 0$ .

We now define the **global** inner product of  $\phi, \psi$  as

$$(\phi, \psi) = \int_M \langle \phi, \psi \rangle \eta,$$

where  $\eta$  is the volume form of  $M$ .

- **The exterior differential operator  $d$  and its co-operator** Denote by  $\Lambda^r(M)$  the set of smooth  $r$ -forms on  $M$ . Let  $(\cdot, \cdot)$  be the (global) inner product defined above. As the formal adjoint operator of the exterior differential operator  $d$ , the *codifferential operator*  $\delta : \Lambda^{r+1}(M) \rightarrow \Lambda^r(M)$  is defined by, for every  $\phi \in \Lambda^r(M), \psi \in \Lambda^{r+1}(M)$ ,

$$(d\phi, \psi) = (\phi, \delta\psi).$$

- **Hodge-star operator.** In order to find the expression of the codifferential operator  $\delta$ , we introduce the Hodge-star operator  $*$ , which is an isomorphism  $*$  :  $\Lambda^r(M) \rightarrow \Lambda^{m-r}(M)$  defined by, for every  $\phi, \eta \in \Lambda^r(M)$ ,

$$\phi \wedge (*\psi) = \langle \phi, \psi \rangle \eta.$$

Let  $\omega$  be a  $r$ -form. Let  $(U, x^i)$  be a local coordinate. We write

$$\omega|_U = \frac{1}{r!} \sum_{i_1, \dots, i_r} a_{i_1 \dots i_r} dx^{i_1} \wedge \dots \wedge dx^{i_r}.$$

Then

$$*\omega = \frac{\sqrt{G}}{r!(m-r)!} \delta_{i_1 \dots i_m}^{1 \dots m} a^{i_1 \dots i_r} dx^{i_{r+1}} \wedge \dots \wedge dx^{i_m},$$

where

$$a^{i_1 \dots i_r} = g^{i_1 j_1} \dots g^{i_r j_r} a_{j_1 \dots j_r},$$

and  $\delta_{i_1 \dots i_m}^{1 \dots m}$  is the Levi-Civita permutation symbol, i.e.  $\delta_{i_1 \dots i_m}^{1 \dots m} = 1$  if  $(i_1 \dots i_m)$  is an even permutation of  $(1 \dots m)$ ,  $\delta_{i_1 \dots i_m}^{1 \dots m} = -1$  if  $(i_1 \dots i_m)$  is an odd permutation of  $(1 \dots m)$ ,  $\delta_{i_1 \dots i_m}^{1 \dots m} = 0$  otherwise. It can be shown that  $*\omega$  is independent of the choice of local coordinates. So  $*\omega$  is a globally defined  $(m-r)$ -form (it can be regarded as an alternative definition). The operator  $*$  which sends  $r$ -forms to  $(m-r)$ -forms.

It has the following properties, for any  $r$ -forms  $\phi$  and  $\psi$ :

$$(1) \phi \wedge *\psi = \langle \phi, \psi \rangle \eta,$$

$$(2) *\eta = 1, *1 = \eta,$$

$$(3) *(*\phi) = (-1)^{r(m+1)} \phi,$$

$$(4) (*\phi, *\psi) = (\phi, \psi).$$

- **Expression of the codifferential operator  $\delta$  in terms of the Hodge-Star operator.** Define  $\delta = (-1)^{mr+1} * \circ d \circ * : \Lambda^{r+1}(M) \rightarrow \Lambda^r(M)$ , where  $\Lambda^r(M)$  is the set of smooth  $r$ -forms, is called the *codifferential operator*. It is easy to verify that  $\delta \circ \delta = 0$ . We also have the following very important property for  $\delta$ : For  $\phi \in \Lambda^r(M)$ ,  $\psi \in \Lambda^{r+1}(M)$ , we have

$$(d\phi, \psi) = (\phi, \delta\psi),$$

i.e.  $\delta$  is conjugate to  $d$ . So  $(-1)^{mr+1} * \circ d \circ *$  is the expression of the codifferential operator  $\delta$ .

*Proof.* Note

$$\begin{aligned}
d(\phi \wedge * \psi) &= d\phi \wedge * \psi + (-1)^r \phi \wedge d(* \psi) \\
&= d\phi \wedge * \psi + (-1)^r (-1)^{mr+r} \phi \wedge *(d * \psi) \\
&= d\phi \wedge * \psi - \phi \wedge * d\psi.
\end{aligned}$$

Then desired identity is obtained by applying the Stokes theorem.

- **Hodge-Laplace operator.** We define the Hodge-Laplace operator

$$\tilde{\Delta} = d\delta + \delta d : \Lambda^r(M) \rightarrow \Lambda^r(M).$$

For  $f \in C^\infty(M)$ , then  $\delta(f) = 0$ , so

$$\tilde{\Delta}(f) = \delta(df) = - * d * df, \quad \tilde{\Delta} f \eta = * \tilde{\Delta} f = -d * df.$$

Let  $(U, x^i)$  be a local coordinate, then

$$df|_U = \frac{\partial f}{\partial x^i} dx^i,$$

$$\begin{aligned}
*df|_U &= \frac{\sqrt{G}}{(m-1)!} \delta_{i_1 \dots i_m}^{1 \dots m} g^{i_1 j} \frac{\partial f}{\partial x^j} dx^{i_2} \wedge \dots \wedge dx^{i_m} \\
&= \sqrt{G} \sum_{i=1}^m (-1)^{i+1} g^{ij} \frac{\partial f}{\partial x^j} dx^1 \wedge \dots \wedge \hat{dx}^i \wedge \dots \wedge dx^m.
\end{aligned}$$

Hence

$$\begin{aligned}
(\tilde{\Delta} f) \eta|_U &= -d(*df)|_U = -\frac{\partial}{\partial x^i} \left( \sqrt{G} g^{ij} \frac{\partial f}{\partial x^j} \right) dx^1 \wedge \dots \wedge dx^m \\
&= -\Delta f \eta|_U.
\end{aligned}$$

This tells us

$$\tilde{\Delta} f = -\Delta f.$$

So  $-\tilde{\Delta}$  when acts on  $C^\infty(M)$  is the Beltrami-Laplace operator  $\Delta$ .

- **Hodge Theory.** In this section, we denote the Hodge-Laplace operator by  $\Delta$ . Let  $\mathcal{H}^r(M) = \ker \Delta$  and  $\mathcal{H} = \bigoplus \mathcal{H}^r(M)$ . Let  $\Lambda^*(M) = \bigoplus_{r=0}^{\infty} \Lambda^r(M)$ .

**The Hodge theorem** *Let  $(M, g)$  be an  $n$ -dimensional compact oriented Riemannian manifold without boundary. For each integer  $0 \leq r \leq n$ ,  $\mathcal{H}^r(M)$  is finite dimensional, and there exists a bounded linear operator  $G : \Lambda^*(M) \rightarrow \Lambda^*(M)$  (called Green's operator) such that*

- (a)  $\ker G = \mathcal{H}$ ;
- (b)  $G$  keeps types, and commute with the operators  $*$ ,  $d$  and  $\delta$ ;
- (c)  $G$  is a compact operator, i.e. the closure of image of an arbitrary bounded subset of  $\Lambda^*(M)$  under  $G$  is compact;
- (d)  $I = \mathcal{H} + \Delta \circ G$ , where  $I$  is the identity operator, and  $\mathcal{H}$  is the orthogonal projection from  $\Lambda^*(M)$  to  $\mathcal{H}$  with respect to the inner product  $(\cdot, \cdot)$ .

From the Hodge theorem, since  $I = \mathcal{H} + \Delta \circ G$ , we can write (called the Hodge-decomposition)

**Corollary( Hodge-decomposition)**

$$\begin{aligned}
 \Lambda^r(M) &= \Delta(\Lambda^r(M)) \oplus \mathcal{H}^r(M) \\
 &= d\delta\Lambda^r(M) \oplus \delta d\Lambda^r(M) \oplus \mathcal{H}^r(M) \\
 &= d\Lambda^{r-1}(M) \oplus \delta\Lambda^{r+1}(M) \oplus \mathcal{H}^r(M).
 \end{aligned}$$